

calculations we try to avoid specifying such a basis until we need an objectification first by γ_0 then some extension direction. Anyway, the STA idea is founded on the orthonormal $\{\gamma_0, \gamma_1, \gamma_2, \gamma_3\}$.

We see that the Minkowski \mathcal{B} -bivector for the line-extension-like information (7.5) $\sigma_k \equiv \gamma_k \wedge \gamma_0$ has (space) Descartes *extension parity inversion* $\overline{\sigma_k} = -\sigma_k \equiv \overline{\gamma_k \wedge \gamma_0} \in \mathcal{G}_{1,3}(\mathbb{R})$, and $\overline{\sigma_k} = -\sigma_k$. While the angular-area-extension-like information bivector (7.6) $i_k \equiv \gamma_i \wedge \gamma_j$ is *parity inversion invariant* $\overline{i_k} = i_k \equiv \overline{\gamma_i \wedge \gamma_j} = \gamma_i \wedge \gamma_j$, of course, not reversion invariant $\widetilde{i_k} = -i_k \equiv \overline{\gamma_j \wedge \gamma_i}$.⁴¹⁰ Here it is urgent to note $\overline{\gamma_k \wedge \gamma_0} = -\gamma_k \wedge \gamma_0 = \gamma^k \wedge \gamma^0$.

In general, we write the reciprocal bivector basis as $\gamma^{\mu\nu} = \gamma^\mu \gamma^\nu = \gamma^\mu \wedge \gamma^\nu$, for $\nu \neq \mu = 0, 1, 2, 3$

$$(7.43) \quad \{\gamma^{10}, \gamma^{20}, \gamma^{30}, \gamma^{23}, \gamma^{31}, \gamma^{12}\} = \{\gamma_{01}, \gamma_{02}, \gamma_{03}, \gamma_{23}, \gamma_{31}, \gamma_{12}\} = \{\overline{\gamma_\mu \wedge \gamma_\nu}\} \\ = \{\gamma_{01}, \gamma_{02}, \gamma_{03}, \gamma_{23}, \gamma_{31}, \gamma_{12}\} = \{-\gamma_{10}, -\gamma_{20}, -\gamma_{30}, -\gamma_{32}, -\gamma_{13}, -\gamma_{21}\}$$

The trivector basis is just $\{i\gamma_\mu\}$ and its reciprocal $\{i\gamma^\mu\}$ due to (7.15) and (7.40).

7.1.3.6. The STA Bivector Field F of STA in the Information Development \mathfrak{D} -space of Physics

We recall the stationary Pauli Basis (6.119) for the geometric algebra $\mathcal{G}_3(\mathbb{R})$ for 3-space

$$(7.44) \quad \{1, \sigma_1, \sigma_2, \sigma_3, i_1 = i\sigma_1 = \sigma_3\sigma_2, i_2 = i\sigma_2 = \sigma_1\sigma_3, i_3 = i\sigma_3 = \sigma_2\sigma_1, i := \sigma_3\sigma_2\sigma_1\}, \gamma_0 \gamma_0$$

Right multiplying this with the information development measure unit γ_0 we get the odd basis $\mathcal{G}_{1,3}^-(\mathbb{R}) \subset \mathcal{G}_{1,3}(\mathbb{R})$ for STA in \mathfrak{D} -space concerning 1-vectors and trivectors

$$(7.45) \quad \{\gamma_0, \gamma_1, \gamma_2, \gamma_3, \gamma_3\gamma_0\gamma_2, \gamma_1\gamma_0\gamma_3, \gamma_2\gamma_0\gamma_1, -\gamma_3\gamma_2\gamma_1\} \cdot \gamma_0$$

Right multiplying once again with the unit measure γ_0 we get an even basis $\mathcal{G}_{1,3}^+(\mathbb{R})$ for STA

$$(7.46) \quad \{1, \gamma_1\gamma_0, \gamma_2\gamma_0, \gamma_3\gamma_0, \gamma_2\gamma_3, \gamma_3\gamma_1, \gamma_1\gamma_2, \gamma_0\gamma_3\gamma_2\gamma_1\}$$

To explicit this six element basis with bivectors in the even algebra $\gamma_0 \subset \mathcal{G}_{1,3}(\mathbb{R})$ we write

$$(7.47) \quad \{1, \gamma_{10} = \gamma_1 \wedge \gamma_0, \gamma_{20} = \gamma_2 \wedge \gamma_0, \gamma_{30} = \gamma_3 \wedge \gamma_0, \gamma_{23} = \gamma_2 \wedge \gamma_3, \gamma_{31} = \gamma_3 \wedge \gamma_1, \gamma_{12} = \gamma_1 \wedge \gamma_2, i\}$$

We note that $\gamma_{\mu\nu} = \gamma_\mu \gamma_\nu = \gamma_\mu \wedge \gamma_\nu$, for $\nu \neq \mu$ are the bivector basis $\{\gamma_\mu \wedge \gamma_\nu\}$ for $F = \langle X \rangle_2 \in \mathcal{G}_{1,3}(\mathbb{R})$.

For a general bivector field $F \in \mathcal{G}_{1,3}(\mathbb{R})$, we write the expansion

$$(7.48) \quad F = \frac{1}{2} F^{\mu\nu} (\gamma_\mu \wedge \gamma_\nu) = \frac{1}{2} F^{\mu\nu} \gamma_\mu \gamma_\nu = \frac{1}{2} F^{\mu\nu} \gamma_{\mu\nu} \in \mathcal{G}_{1,3}^+(\mathbb{R}),$$

with six real scalar coordinate *components*

$$(7.49) \quad F^{\mu\nu} = (\gamma^\nu \wedge \gamma^\mu) \cdot F = (\gamma^\nu \gamma^\mu) \cdot F = \gamma^{\nu\mu} \cdot F \in \mathbb{R}, \quad \text{as [24]p.2(13),}$$

with reason, coordinates after 1-vector *directions*: first γ^μ , then γ^ν , as $\gamma^\nu \cdot (\gamma^\mu \cdot F) = (\gamma^\nu \wedge \gamma^\mu) \cdot F$.

In the matrix tradition $F^{\mu\nu}$ form an 4×4 antisymmetric contravariant matrix as a rank-2 tensor.

In Space-Time Algebra it is a bivector field that can be resolved in six independent components.

The $\frac{1}{2}$ in (7.48) comes from the anti-symmetry in the sum $\frac{1}{2} \sum_{\mu, \nu=0}^3 F^{\mu\nu} \gamma_{\mu\nu}$, (12 terms and 4 zeros).

7.1.3.7. Difference Between the Pseudoscalar Concepts for \mathfrak{D} -space and 3-space

– The \mathfrak{D} -space concept we endow with the geometric algebra $\mathcal{G}_{1,3}(\mathbb{R})$ managed by the four-vector top *grade four direction* helicity pseudoscalar concept unit $i := \gamma_1 \gamma_2 \gamma_3 \gamma_0 = \gamma_0 \gamma_3 \gamma_2 \gamma_1$, which is invariant to *Clifford reversion conjugation* $\widetilde{i} = i$, and commute with every even multivector

$$(7.50) \quad X_+ i = i X_+, \quad \text{of the form } X_+ = \langle X \rangle_0 + \langle X \rangle_2 + \langle X \rangle_4 \in \mathcal{G}_{1,3}^+(\mathbb{R}) \subset \mathcal{G}_{1,3}(\mathbb{R}),$$

but anticommute with odd elements of this algebra

$$(7.51) \quad i X_- = -X_- i, \quad \text{of the form } X_- = \langle X \rangle_1 + \langle X \rangle_3 \in \mathcal{G}_{1,3}^-(\mathbb{R}) \subset \mathcal{G}_{1,3}(\mathbb{R}).$$

⁴¹⁰ Were $k \neq j \neq i \neq k$, all $\in \{3, 2, 1\}$ are permutations.

– Opposite in 3-space endowed by $\mathcal{G}_3(\mathbb{R})$ managed by the trivector of top *grade three* chiral volume pseudoscalar concept unit $i := \sigma_3 \sigma_2 \sigma_1$ *direction*, that reverses as $i^\dagger = -i$, and commutes with all elements $A \in \mathcal{G}_3(\mathbb{R}) : iA = Ai$. Yet invariant by *Clifford conjugation* $\widetilde{i} = i \in \mathcal{G}_3(\mathbb{R})$.

7.1.4. The Odd and Even Part of the Geometric Space-Time Algebra $\mathcal{G}_{1,3}(\mathbb{R})$

The general multivector form (7.27)–(7.33) can be decomposed into a sum of *even* and *odd* parts

$$(7.52) \quad X = X_+ + X_- = \langle X \rangle_+ + \langle X \rangle_-, \quad \text{where}$$

$$(7.53) \quad \langle X \rangle_+ = \langle X \rangle_0 + \langle X \rangle_2 + \langle X \rangle_4 = X_+ = \alpha + F + vi \in \mathcal{G}_{1,3}^+(\mathbb{R}) \subset \mathcal{G}_{1,3}(\mathbb{R}),$$

$$(7.54) \quad \langle X \rangle_- = \langle X \rangle_1 + \langle X \rangle_3 = X_- = x + iy \in \mathcal{G}_{1,3}^-(\mathbb{R}) \subset \mathcal{G}_{1,3}(\mathbb{R}).$$

This decomposed concept idea we write as $\mathcal{G}_{1,3}(\mathbb{R}) = \mathcal{G}_{1,3}^+(\mathbb{R}) \oplus \mathcal{G}_{1,3}^-(\mathbb{R})$.

Algebraically this can be formed by $X_\pm = \frac{1}{2}(X \mp iXi)$, an idea taken from Hestenes [23](26)p.6.

7.1.4.2. The Transcendental Ignorance of the Odd part of the Geometric Algebra for \mathfrak{D} -space

The founding *substance* we have introduced as (essential ideal) three *subjects* of Descartes' extensional unit *directions* $\gamma_1, \gamma_2, \gamma_3$, that in their idea are *eternal* unchangeable orthonormal autonomous for any *entity* in Nature. The extra local cyclic unit count with (essential ideal) *direction* γ_0 into the *FUTURE* is the idea of a *persistent* chronometric *subject* for development.

The founding basis $\{\gamma_0, \gamma_1, \gamma_2, \gamma_3\}$ for the $\mathcal{G}_{1,3}(\mathbb{R})$ is the concept idea of *subjects* of the concept of a \mathfrak{D} -space *substance* of physics. The linear construction (7.10), (7.30) $x = x^\mu \gamma_\mu$ 1-vector and its algebraic dual $ix = x^\mu i\gamma_\mu$ or (7.32) $iy = y^\mu i\gamma_\mu$ trivector is a pure substantial mathematical construction that has no direct object image for the intuition of natural space for physics.

The odd standard basis $\{\gamma_0, \gamma_1, \gamma_2, \gamma_3\}$ is the *directional* subject foundation for our algebra $\mathcal{G}_{1,3}(\mathbb{R})$ established from a traditional non-Euclidean linear additional vector space $V_{1,3}(\mathbb{R}) \sim \mathbb{R}^4$ by (7.10). Mutual multiplication of the odd elements $X_- \in \mathcal{G}_{1,3}^-(\mathbb{R})$ are lifted into the even algebra $\mathcal{G}_{1,3}^+(\mathbb{R})$ and further multiplication will stay closed inside this even algebra. (Out of the transcendental odd?) Now we will ignore the odd case and concentrate on the even closed algebra $\mathcal{G}_{1,3}^+(\mathbb{R})$.

7.1.4.3. The Even Closed Geometric Algebra of \mathfrak{D} -space

The even closed algebra $\mathcal{G}_{1,3}^+(\mathbb{R})$ consists of scalars

$$\langle X \rangle_0 = \alpha \in \mathbb{R}, \quad (7.29)$$

and pseudoscalars

$$\langle X \rangle_4 = vi, \quad (7.33)$$

together with the bivector field concept in \mathfrak{D} -space of physics

$$\langle X \rangle_2 = F = (F^{10} \gamma_1 \gamma_0 + F^{20} \gamma_2 \gamma_0 + F^{30} \gamma_3 \gamma_0) + (F^{32} \gamma_2 \gamma_3 + F^{13} \gamma_3 \gamma_1 + F^{21} \gamma_1 \gamma_2) \quad (7.31) \Leftrightarrow$$

$$(7.55) \quad \langle X \rangle_2 = F = (F^{10} \mathcal{B}_1 + F^{20} \mathcal{B}_2 + F^{30} \mathcal{B}_3) + (F^{32} i\mathcal{B}_1 + F^{13} i\mathcal{B}_2 + F^{21} i\mathcal{B}_3),$$

that is (a Minkowski \mathcal{B} -plane bivector) + (a Cartesian extension plane bivector). We recall that these two bivector parts stand in dual *quality* to each other. We have the field bivector

$$(7.56) \quad F = \sum_{k=1}^3 F^{k0} \mathcal{B}_k + \frac{1}{2} \sum_{j,k=1}^3 F^{kj} \gamma_k \gamma_j \cong F^{k0} \sigma_k + \frac{1}{2} F^{kj} \sigma_j \sigma_k = \vec{E} + i\vec{B}.$$

The left side sum is the bivector field in $\mathcal{G}_{1,3}^+(\mathbb{R})$ that mimic the electromagnetic tensor field.

The right side is e.g. the electric field 1-vector plus the magnetic field bivector in $\mathcal{G}_3(\mathbb{R})$.

The reader is urged to consider the connection idea $\mathcal{G}_{1,3}^+(\mathbb{R}) \sim \mathcal{G}_3(\mathbb{R})$ as a *subject~object* principle.

7.1.4.4. The Even Geometric Spinor Quality in the Algebra $\mathcal{G}_{1,3}^+(\mathbb{R})$ of \mathfrak{D} -space

The general even multivector element form (7.28) is then written as a *spinor* (wavefunction)

$$(7.57) \quad \left. \begin{aligned} \psi &= \alpha + F + vi \\ \widetilde{\psi} &= \alpha - F + vi \end{aligned} \right\} \in \mathcal{G}_{1,3}^+(\mathbb{R}) \subset \mathcal{G}_{1,3}(\mathbb{R}).$$