

6.3. The 3-space Structure *Quality* Described by Multivectors

A general arbitrary 3-multivector is constructed from four *primary qualities of grades* $pqg-0 + pqg-1 + pqg-2 + pqg-3$. The approach tries to add these different qualities together. A classical space solid object we have the tradition to define in a Euclidean vector space $(V_3, \mathbb{R}) \sim \mathbb{R}_1^3$ fixed to a standard *orthonormal dextral basis* $\{\sigma_1, \sigma_2, \sigma_3\}$ of three 1-vector objects. All 1-vectors in this Euclidean vector space possess the *primary qualities of first grade (pqg-1)*. This basis implies the transversal bivector basis $\{i_1, i_2, i_3\} = \{\sigma_3\sigma_2, \sigma_1\sigma_3, \sigma_2\sigma_1\}$, (6.31), (6.34) as a generator of all planes possessing *pqg-2 quality*. On top of that we have trivector chiral pseudoscalar unit $i := \sigma_3\sigma_2\sigma_1$ generating the chiral volume possessing *pqg-3 quality*.

In all, we expand a multivector form from the geometric algebra $\mathcal{G}_3 = \mathcal{G}_3(\mathbb{R}) = \mathcal{G}(V_3, \mathbb{R})$ as

$$(6.59) \quad A = \alpha + x_1\sigma_1 + x_2\sigma_2 + x_3\sigma_3 + \beta_1 i_1 + \beta_2 i_2 + \beta_3 i_3 + v i$$

$$A = \langle A \rangle_0 + \langle A \rangle_1 + \langle A \rangle_2 + \langle A \rangle_3$$

Refer to: (5.59) | (6.29) | (6.32), (6.33), (6.36) | (6.15) in 6.2.2

We have separated the multivector concept into four different *primary quality grades*

$$(6.60) \quad \langle A \rangle_0 = \alpha, \quad pqg-0, \text{ scalar, } \dim(\mathbb{R}) = 1,$$

$$(6.61) \quad \langle A \rangle_1 = \mathbf{a} = x_1\sigma_1 + x_2\sigma_2 + x_3\sigma_3, \quad pqg-1, \text{ 1-vector, } \dim(V_3) = 3,$$

$$(6.62) \quad \langle A \rangle_2 = \mathbf{b}i = (\beta_1\sigma_1 + \beta_2\sigma_2 + \beta_3\sigma_3)i = \beta_1 i_1 + \beta_2 i_2 + \beta_3 i_3, \quad pqg-2, \text{ bivector, } \dim(V_3) = 3,$$

$$(6.63) \quad \langle A \rangle_3 = v i, \quad \text{chiral volume pseudoscalar, } pqg-3, \text{ trivector, } \dim(\mathbb{R}) = 1.$$

The multivector idea over a geometric product algebra has a linear addition structure of multiple *primary qualities of grades* we call it $\mathcal{G}_3(\mathbb{R}) = \mathcal{G}(V_3, \mathbb{R})$ over a Euclidean vector space V_3 , with which we try to describe the *local topological form structure* of the physical 3-space.

This linear addition structure for multivectors over a real field \mathbb{R} we express

$$(6.64) \quad A = \alpha + \mathbf{a} + \mathbf{b}i + v i.$$

The linear algebra $\mathcal{G}_3 = \mathcal{G}_3(\mathbb{R})$ has more dimension from the different *grade qualities*, in all

$$(6.65) \quad \dim(\mathcal{G}_3) = 1 + 3 + 3 + 1 = 8. \quad \text{In generell, } \dim(\mathcal{G}_n) = \sum_{r=0}^n \binom{n}{r} = 2^n$$

The mixed basis for the whole linear algebra \mathcal{G}_3 is $\{1, \sigma_1, \sigma_2, \sigma_3, \sigma_3\sigma_2, \sigma_1\sigma_3, \sigma_2\sigma_1, \sigma_3\sigma_2\sigma_1\}$.

6.3.2. The Even and the Odd Geometric Algebra

When we split this algebra in even \mathcal{G}_3^+ and odd \mathcal{G}_3^- , so $\mathcal{G}_3 = \mathcal{G}_3^- + \mathcal{G}_3^+$, we write the multivector

$$(6.66) \quad A = \langle A \rangle^+ + \langle A \rangle^- = \langle A \rangle_{0,2}^+ + \langle A \rangle_{1,3}^- = \underbrace{\alpha + \mathbf{b}i}_+ + \underbrace{\mathbf{a} + v i}_-$$

where the even algebra \mathcal{G}_3^+ is a *closed* multiplication spinor algebra, while \mathcal{G}_3^- is open, not closed. The spinor multivector subject can be expressed as a 1-rotor multiplied with a real dilation factor

$$(6.67) \quad \langle A \rangle^+ = \langle A \rangle_{0,2}^+ = \langle A \rangle_0 + \langle A \rangle_2 = \alpha + \mathbf{b}i = \mathbf{r} \cdot \mathbf{u} + \mathbf{r} \wedge \mathbf{u} = \mathbf{r} \mathbf{u} = \rho U = \rho e^{i\theta} \in \mathcal{G}_3^+(\mathbb{R}),$$

as in (5.163) origin in (5.97), where the rotation is in the plane of the bivector $\mathbf{b}i = \mathbf{r} \wedge \mathbf{u}$.

You may define two object 1-vectors whose product spinor $\mathbf{r} \mathbf{u}$ determines the rotation, and dilation $\rho = |\mathbf{r} \mathbf{u}|$. (We prefer to choose $|\mathbf{u}|=1$ so that $\rho = |\mathbf{r}|$ if it is possible.)²⁹⁷

If you determine the 1-vector $\mathbf{b} = -i(\mathbf{r} \wedge \mathbf{u})$ as an object it will represent the rotation axis for $\langle A \rangle^+$. This rotation axis has a transversal plane *direction* in which \mathbf{r} and \mathbf{u} exist.

The odd algebra $\mathcal{G}_3^-(\mathbb{R})$ part substance of 3-space gives us subjects $\langle A \rangle_{1,3}^- = \mathbf{a} + v i$.

The geometric 1-vector can give us a straight rectilinear translation as in §4.4.2.13 and § 5.3.7.2 by the subject $\mathbf{t} = \mathbf{a}$, or by the 1-vector object \mathbf{a} representing a *pqg-1 direction* and a straight line-segment magnitude $|\mathbf{a}|$ for a physical *entity*. The new *pqg-3* subject $v i$ gives us a *chiral-directed* volume of a solid spatial *entity* object, as a *chiral quality pqg-3*.

²⁹⁷ I think that there could be a little problem here; if we find \mathbf{r} and \mathbf{u} colinear $\mathbf{r} \wedge \mathbf{u} = 0 \Leftrightarrow \mathbf{b} = 0$ then the spinor is a pure scalar $\mathbf{r} \cdot \mathbf{u}$, and by that, does a pure scalar spinor gives any sense, even if it has a final magnitude $|\mathbf{r} \mathbf{u}| > 0$? Oscillations will resolve this.