

5.7.2.3. Minkowski Space with Display of Three Extension Dimensions

We follow the display with the development *direction* γ_0 vertical and the *null directions* turned $\pm 45^\circ$ from this in the \mathcal{B}_1 -bivector-plane. In Figure 5.54 we display the projection of the extension *directions* γ_1 and γ_2 with the hyperbola transformation parameter λ close to one in the defining formula (5.319) as $\gamma_1(\lambda \approx 1)$ and $\gamma_2(\lambda \approx 1)$. The same we do in Figure 5.55 where we further introduce the third extension *direction* γ_3 , with $\lambda \rightarrow 0$; projection of $\gamma_3(\lambda = 0)$ displayed *antiparallel* \Downarrow to γ_0 , $\gamma_3 \parallel \gamma_0$. Chosen with the opposite orientation of the development *direction*. But attention; $\gamma_0 \cdot \gamma_3 = 0$; there are all algebraically *orthogonal* $\gamma_0 \cdot \gamma_k = 0$, for $k=1,2,3$, (5.303), and we too presume $\gamma_1 \cdot \gamma_2 = \gamma_2 \cdot \gamma_3 = \gamma_3 \cdot \gamma_1 = 0$.

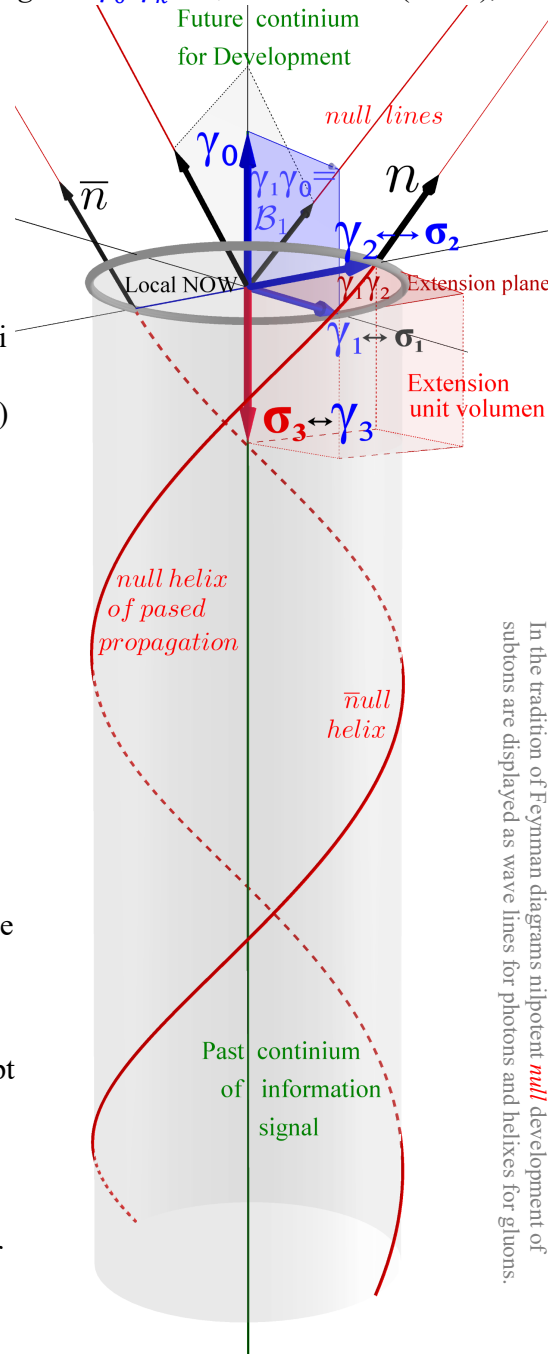
The projection mapping of the Minkowski basis $\{\gamma_0, \gamma_1, \gamma_2, \gamma_3\}$ with signatures $(+, -, -, -)$ into the Euclidean basis $\{\sigma_1, \sigma_2, \sigma_3\}$ with signatures $(+, +, +)$ is performed by the multiplication operations (5.328)

$$(5.341) \quad \begin{aligned} \gamma_1 &\rightarrow \sigma_1 \equiv \gamma_1 \gamma_0, & \sigma_1 &\rightarrow \gamma_1 \equiv \sigma_1 \gamma_0 \\ \gamma_2 &\rightarrow \sigma_2 \equiv \gamma_2 \gamma_0, & \text{convers} & \sigma_2 &\rightarrow \gamma_2 \equiv \sigma_2 \gamma_0 \\ \gamma_3 &\rightarrow \sigma_3 \equiv \gamma_3 \gamma_0, & \sigma_3 &\rightarrow \gamma_3 \equiv \sigma_3 \gamma_0 \end{aligned}$$

In Figure 5.55 we will imagine that the third Minkowski \mathcal{B}_3 -bivector-plane $\mathcal{B}_3^2=1$, $\mathcal{B}_3 := \gamma_3 \gamma_0$ has the same start *direction* as $\mathcal{B}_1 := \gamma_1 \gamma_0 \rightarrow \sigma_1$. (\mathcal{B}_3 as \mathcal{B}_2 is un-displayed) The invariant nilpotent *null directions* n and \bar{n} in the \mathcal{B}_1 -bivector-plane is a translation invariant moved one unit along the $\pm \gamma_2$ *direction* of the Figure 5.55 display. The idea is that the *entity* represented by this frame is rotating with an oscillation in the *extension plane* $\gamma_1 \gamma_2$ $\circ = e^{\gamma_1 \gamma_2 \theta} = e^{i\theta}$, relative to an external laboratory. By this oscillating rotation of both \mathcal{B}_3 and \mathcal{B}_1 -bivector *directions* the nilpotent isotropic *directions* are twisted up in two opposite equal-orientated *null helixes* stretched out in the past as the oscillator is propagating into the future. It is the nilpotence of the *null basis* $\{n, \bar{n}\}$ (5.310) $n^2 = \bar{n}^2 = 0$ (5.311) that makes this invariant stretching possible carrying an information signal through the development from the past we call the *subton memory space of extension*.

This signal of information through space is what we in chapter I. 3.4 called a subton, which was displayed in Figure 3.13. where we did not have the nilpotent concept idea, but implicit had, that carrier time always is $t_{\bar{c}} = 0$ throughout the propagation as described in section I. 3.5.2, and which transfer of information in principle is displayed in Figure 3.14. The subton idea is the a priori foundation concept of a $U(1)$ spin-1-rotor oscillator propagating with the speed of information, which we set isometric equal to one, where the autonomous measure is one count of *quantum* in the oscillation of the subton *entity* development. The primitive subton idea is the foundation for the concept of a photon (and properly also gluons).

Figure 5.55 The four-dimensional Minkowski space projected into the figure plane. The extension space is generated by a plane cyclic oscillating rotation producing a helix-formed nilpotent *null* curve of development. The frame $\{\sigma_k\}$ performs a rotating oscillation around



In the tradition of Feynman diagrams nilpotent *null* development of subtons are displayed as wave lines for photons and helixes for gluons.

Geometric Critique on the a priori of Physics

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5.7.3. The Paravector Space and the Minkowski 1-vector in STA

In § 5.5.4.2 (5.240) we defined the foundation of a paravector as a scalar plus a 1-vector

$$(5.342) \quad p = \lambda_0 + \mathbf{p} = \lambda_0 1 + \lambda_1 \mathbf{u}.$$

The 1-vector can have different directions $\mathbf{p} = \lambda_1 \sigma_1 + \lambda_2 \sigma_2$ in a Euclidean plane $\{\sigma_1, \sigma_2\}$ or in space $\{\sigma_k\}$, $\mathbf{p} = \lambda_k \sigma_k$, for $k=1,2,3$.²⁷⁸ The para-multi-vectors of *grade* form $\langle A \rangle_0 + \langle A \rangle_1$ are then

$$(5.343) \quad p = \lambda_0 + \mathbf{p} = \lambda_0 1 + \lambda_k \sigma_k.$$

The pure magnitude of the paravector is given from (5.253)–(5.255) as

$$(5.344) \quad |p| = |\bar{p}| = \sqrt{|p\bar{p}|} = \sqrt{|\lambda_0^2 - \lambda_1^2 - \lambda_2^2 - \lambda_3^2|} \geq 0.$$

The paravector space is supported from the mixed basis $\{1, \sigma_1, \sigma_2, \sigma_3\}$ of form $\langle A \rangle_0 + \langle A \rangle_1$. To see how the paravector concept looks like in Minkowski space we map by multiplication operation from the right with the *quantum count* of development expressed as the 1-vector unit γ_0

$$(5.345) \quad p = p\gamma_0 = \lambda_0 \gamma_0 + \mathbf{p}\gamma_0 = \lambda_0 \gamma_0 + \lambda_k \sigma_k \gamma_0 = \lambda_0 \gamma_0 + \lambda_k \gamma_k = \lambda_\mu \gamma_\mu. \quad \text{Use (5.328) } \sigma_k \gamma_0 = \gamma_k.$$

This is a 1-vector in a four-dimensional Minkowski space supported as a span from the orthonormal 1-vector basis $\{\gamma_\mu\}$, for $\mu = 0,1,2,3$.²⁷⁹ Here the multivector form is a pure $\langle A \rangle_1$, i.e., a *primary quality of first grade (pqg-1)* with the *development count* as the *quantity* unit. The geometric algebra founded on this basis $\{\gamma_\mu\}$ we as Hestenes [6] 1966 call Space-Time-Algebra (STA). The 1-vector basis $\{\gamma_\mu\}$ is sometimes called a Dirac basis of STA -or just- a Minkowski basis. The magnitude of such a 1-vector $p = \lambda_\mu \gamma_\mu$ in the Minkowski metric is given by the quadratic form

$$(5.346) \quad p^2 = (\lambda_\mu \gamma_\mu)^2 = \lambda_0^2 - \lambda_1^2 - \lambda_2^2 - \lambda_3^2 \Rightarrow |p| = \sqrt{|p^2|} = \sqrt{|\lambda_0^2 - \lambda_1^2 - \lambda_2^2 - \lambda_3^2|} \geq 0.$$

The two pictures describe the same structure of a physical *entity* therefore the magnitude $|p| = |\bar{p}|$. Expressed from the supporting 1-vector basis $\{\gamma_\mu\}$ the Cartesian basis $\{\sigma_k\} = \{\sigma_1, \sigma_2, \sigma_3\}$ looks like a Minkowski \mathcal{B} -bivector basis $\{\mathcal{B}_k\}$ by the map (5.326) $\mathcal{B}_k := \gamma_k \gamma_0 \leftrightarrow \sigma_k$

The Cartesian basis $\{\sigma_k\}$ supports the Euclidean space by 1-vectors describing the Cartesian extension of locality for an *entity* in physics. Adding a scalar dimension to the Cartesian frame gives us a paravector (5.343) with a mixed basis $\{1, \mathbf{u}\}$, $\{1, \sigma_1, \sigma_2\}$ or $\{1, \sigma_1, \sigma_2, \sigma_3\}$. This paravector picture split the picture of STA developing space with four *direction* dimensions into a scalar part without *direction* and one, two or three Cartesian extension *directions* σ_k . This agrees with the classical way to interpret space extension with a distance measure separated from the development of time as a pure scalar measure from a clock without extensive *direction*. Contrary if the local cyclic oscillating rotation clock produces the *direction* of *quantum counts* γ_0 and we can right multiply it to the paravector basis using (5.328) and get the STA basis

$$(5.347) \quad \{1\gamma_0, \sigma_1\gamma_0, \sigma_2\gamma_0, \sigma_3\gamma_0\} = \{\gamma_0, \gamma_1, \gamma_2, \gamma_3\},$$

then γ_0 by *inheritance* to the common measure of extensions gives the performed development.

Having a single 1-vector *direction* \mathbf{u} in a Euclidean space $\mathbf{u}^2=1$, we can choose it as a frame basis 1-vector e.g., $\sigma_1 = \mathbf{u}$. We multiply this with $\gamma_0 \gamma_0 = \gamma_0^2 = 1$ and get the mapping operation

$$(5.348) \quad \mathbf{u} = \mathbf{u}\gamma_0\gamma_0 = \sigma_1\gamma_0\gamma_0 = \gamma_1\gamma_0 = \mathcal{B}_1 = \mathcal{B}$$

We hereby see that the Euclidean 1-vector *direction* \mathbf{u} is synonymous with a Minkowski \mathcal{B} -bivector including an autonomous development *quantum* measure count of its own unit extension.

We note that these two concepts have the same Clifford signature (+): $\mathbf{u}^2 = 1$ and $\mathcal{B}^2=1$. Physical *entities* that we represent by Euclidean 1-vectors are in Space-Time Algebra (STA) represented by \mathcal{B} -bivectors with positive signature (+).

²⁷⁸ Here we use Einstein sum convention $\lambda_k \sigma_k = \sum_{k=1}^3 \lambda_k \sigma_k$ over double indices, and low indices for the covariant coordinates in the orthonormal frame of the Euclidian space, to keep intuition simple. More about $k=3, \{\sigma_k\}$ in chapter 6.

²⁷⁹ Note we commonly use Greek letter indices for four-dimensional Minkowski space and Latin indices for Euclidian 1,2,3 support.