

5.4.5. Rotation Inside the Same Plane Direction

Two different normalized 1-vectors  $\mathbf{u}$  and  $\mathbf{v}$  where  $\mathbf{v} \neq \mathbf{u}$ ,  $|\mathbf{u}|=|\mathbf{v}|=1$  form a bivector  $\mathbf{v} \wedge \mathbf{u}$ , which spans a plane  $\lambda_{\mathbf{u}\mathbf{v}}$ . The product of these two also forms a 2-multi-vector  $U = \mathbf{v}\mathbf{u}$ , which works in the rotor plane  $\lambda_{\mathbf{u}\mathbf{v}}$ , which is the plane of Figure 5.46 in which, there too are indicated two reflecting planes  $\gamma_{\perp\mathbf{u}} \perp \mathbf{u}$  and  $\gamma_{\perp\mathbf{v}} \perp \mathbf{v}$  as small-dotted lines for intuition. These are transversal normal planes to each of the 1-vectors  $\mathbf{u}$  and  $\mathbf{v}$ . These two plane objects  $\gamma_{\perp\mathbf{u}} \perp \lambda_{\mathbf{u}\mathbf{v}}$  and  $\gamma_{\perp\mathbf{v}} \perp \lambda_{\mathbf{u}\mathbf{v}}$  intersect each other in one straight line with the **direction**  $\hat{\omega} \perp \lambda_{\mathbf{u}\mathbf{v}}$ , normal to the paper pointing towards the observer, who is seeded in a remote<sup>o</sup> origo and looks at the plane as an intuiting interpreter of the Figure 5.46 plane. We consider any arbitrary 1-vector  $\mathbf{x}$ , in space  $\mathbb{G}$  whose projection on the  $\lambda_{\mathbf{u}\mathbf{v}}$  plane is showed Figure 5.46.

$\mathbf{x}$  can be reflected around  $\mathbf{u}$ , then we get  $\mathbf{x}' = \mathbf{u}\mathbf{x}\mathbf{u}$ , or reflected along  $\mathbf{u}$ , so that  $\mathbf{x}'' = -\mathbf{u}\mathbf{x}\mathbf{u} = \underline{\mathbf{u}}\mathbf{x}$  (dashed), is called an *irregular rotation*. Hereafter  $\mathbf{x}'$  is reflected around  $\mathbf{v}$ , we get  $\mathbf{v}\mathbf{x}'\mathbf{v}$  or  $\mathbf{x}''$  reflected along  $\mathbf{v}$ , we achieve  $-\mathbf{v}\mathbf{x}''\mathbf{v} = \underline{\mathbf{v}}\mathbf{x}''$ , as another *irregular rotation*.

We compose these two reflections to one *regular rotation*  $\underline{\mathbf{R}} = \underline{\mathbf{v}}\underline{\mathbf{u}}$  (by double sandwiching)

$$(5.190) \quad \underline{\mathbf{R}}\mathbf{x} = \mathbf{U}\mathbf{x}\mathbf{U}^\dagger = \mathbf{v}\mathbf{u}\mathbf{x}\mathbf{u}\mathbf{v} = -\mathbf{v}\mathbf{x}''\mathbf{v} = -\mathbf{v}(-\mathbf{u}\mathbf{x}\mathbf{u})\mathbf{v} \quad \text{alternative} \quad = \mathbf{v}\mathbf{x}'\mathbf{v} = \mathbf{v}(\mathbf{u}\mathbf{x}\mathbf{u})\mathbf{v}.$$

We see that the ambiguity of reflection through a 1-vector is eliminated by this doublet composition, where the multi-vector concept  $\mathbf{v}\mathbf{u}$  is the generator of the transformation.

We define a 2-multivector as an operator called a *rotor*

$$(5.191) \quad \mathbf{U} = \mathbf{v}\mathbf{u} = \mathbf{v} \cdot \mathbf{u} + \mathbf{v} \wedge \mathbf{u} = e^{+\frac{1}{2}\theta} = e^{+i\theta}$$

$$(5.192) \quad \mathbf{U}^\dagger = \mathbf{u}\mathbf{v} = \mathbf{v} \cdot \mathbf{u} - \mathbf{v} \wedge \mathbf{u} = e^{-\frac{1}{2}\theta} = e^{-i\theta}$$

Here we introduce the bivector  $\frac{1}{2}\theta = i\theta$  which represents the rotor area with **direction**, as an argument for this 2-multivector exponential function  $e^{\pm\frac{1}{2}\theta}$ .

The unit of plane **direction** is  $\hat{\mathbf{i}} = \hat{\theta}$ , refer to (5.8), (5.90).

Hence, the regular rotation as a linear transformation is

$$(5.193) \quad \underline{\mathbf{R}}\mathbf{x} = \mathbf{U}\mathbf{x}\mathbf{U}^\dagger = e^{\frac{1}{2}\theta}\mathbf{x}e^{-\frac{1}{2}\theta} = e^\theta\mathbf{x} = e^{i\theta}\mathbf{x}e^{-i\theta} = e^{i2\theta}\mathbf{x}$$

The unitary rotor 2-multivector is rotational invariant in its own plane. This means, its *pqg-1 directions* by the individual 1-vectors  $\mathbf{u}$  and  $\mathbf{v}$  lose their specific meaning, and instead their relative **direction** area  $\frac{1}{2}\theta$ , which is the *pqg-2 direction* of that rotation area.

The 1-rotor is

$$(5.194) \quad \mathbf{U}_\theta = e^{+\frac{1}{2}\theta} = e^{+i\theta},$$

and the *regular rotation* is  $\underline{\mathbf{R}}\mathbf{x} = \mathbf{U}\mathbf{x}\mathbf{U}^\dagger$  in Figure 5.47, (by rotor sandwiching)

If the start 1-vector  $\mathbf{x}$  is in the rotor plane, we can choose  $\mathbf{U} = \mathbf{u}\mathbf{v} = \mathbf{u}\mathbf{x}/|\mathbf{x}|$  the resulting 1-vector  $\mathbf{U}^2\mathbf{x} = e^{i2\theta}\mathbf{x}$  is in that same plane reduced to a reflection (5.170)  $\mathbf{x}' = \mathbf{u}\mathbf{x}\mathbf{u}$  displayed in Figure 5.42 where  $\mathbf{v} \parallel \mathbf{x}$  is omitted. –

The general formulation (5.193)  $\underline{\mathbf{R}}\mathbf{x} = \mathbf{U}\mathbf{x}\mathbf{U}^\dagger$  apply to all vectors in space  $\mathbb{G}$  outside the rotor  $\mathbf{U}$  plane **direction**, see more below § 6.3.3, Figure 6.12.

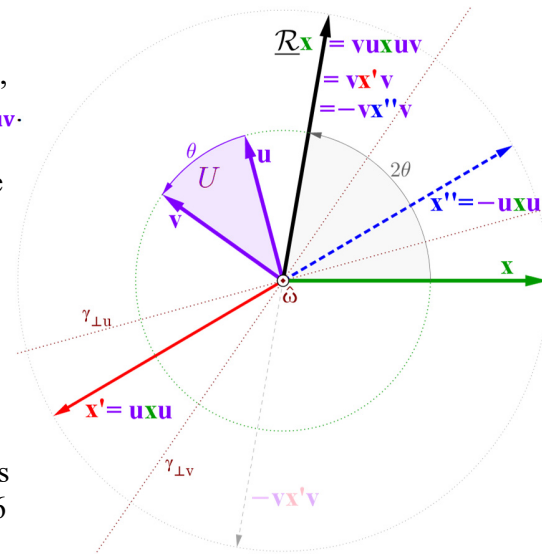


Figure 5.46  $\underline{\mathbf{R}}$ , Rotation in a plane, constituted of two reflections.

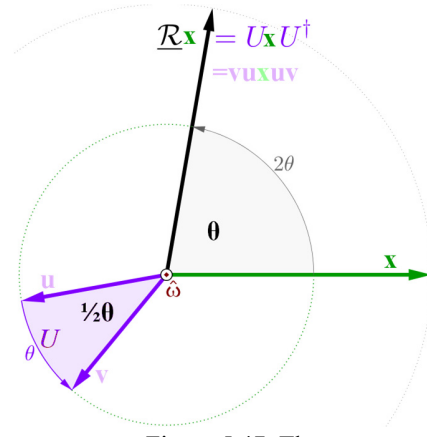


Figure 5.47 The rotation is a linear angular *pqg-2* transformation  $\mathbf{U} = e^{+i\theta}$ .

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This form  $\underline{\mathbf{R}}\mathbf{x} = \mathbf{U}\mathbf{x}\mathbf{U}^\dagger$  is often called the canonical form for any orthogonal transformation  $\underline{\mathbf{R}}$ . This is the *pqg-2* essence of rotation, the **primary quality of the second grade** for all space  $\mathbb{G}$  of physics. A rotation plane **direction** is the same subject all over space  $\mathbb{G}$  (translation invariance). This *pqg-2* symmetry is an a priori foundation for all space  $\mathbb{G}$  in physics.

5.4.5.2. The Half-Angle Rotor of an Euler Rotation

We see that when we have a 1-rotor (5.194)  $\mathbf{U}_\theta = e^{+i\theta}$  the full regular rotation has an argument with the double angle  $\underline{\mathbf{R}}\mathbf{x} = \mathbf{U}\mathbf{x}\mathbf{U}^\dagger = \mathbf{U}^2\mathbf{x} = e^{i2\theta}\mathbf{x}$ . Therefore, we almost always write a 1-rotor with a half angle  $\frac{1}{2}\varphi$  argument causing the full regular rotation with the full angle  $\varphi$  in this

$$(5.195) \quad \mathbf{U}_\varphi = e^{+i\frac{1}{2}\varphi} \Rightarrow \mathbf{U}_\varphi^2 = e^{+i\varphi} \Rightarrow \underline{\mathbf{R}}\mathbf{x} = \mathbf{U}_\varphi\mathbf{x}\mathbf{U}_\varphi^\dagger = e^{+i\frac{1}{2}\varphi}\mathbf{x}e^{-i\frac{1}{2}\varphi} = \mathbf{U}_\varphi^2\mathbf{x} = e^{+i\varphi}\mathbf{x}.$$

This we call an Euler rotation, and if the rotor  $e^{\pm i\frac{1}{2}\varphi}$  is in one fixed plane  $\hat{\mathbf{i}}$  we call it a 1-rotor. The purpose of half angles  $\frac{1}{2}\varphi$  is clear when it comes to Euler rotations in several planes. .

5.4.5.3. The Idea of an Active Rotation

As we know from the treatment in chapters 1, 2 and 3, to experience physics something happens. There always is a development, and the measure for this is a development parameter  $t$  given from a rotating circle oscillator with frequency energy  $\omega$ . We now see that this circle oscillation is a rotation along a Euclidean plane. We simply write this unitary as a plane rotation operator

$$(5.196) \quad e^{\pm i\omega t} = e^{\pm i\varphi} = \mathbf{U}_\varphi^2$$

where  $\hat{\mathbf{i}}$  is the unit bivector **direction** of the Euclidean rotation plane, and  $\omega t = \varphi$  is the active **quantum mechanical** phase angle of the oscillating **entity**, which describes the development.<sup>259</sup> A 1-rotor for an oscillation is often just written  $\mathbf{U}_\varphi = e^{\pm i\frac{1}{2}\varphi} = e^{\pm i\frac{1}{2}\omega t}$ .

5.4.5.4. The Invariant Direction of a Rotor

The rotational invariance of the rotor object  $\mathbf{U} = \mathbf{v}\mathbf{u} = \mathbf{v} \cdot \mathbf{u} + \mathbf{v} \wedge \mathbf{u}$  is illustrated by the difference from Figure 5.46 to Figure 5.47 in accordance with the idea in § 5.2.7.4 and Figure 5.23. By normal duality, we can characterise the rotation plane  $\lambda_{\mathbf{U}} = \lambda_{\mathbf{u}\mathbf{v}} = \lambda_{\mathbf{u}\mathbf{v}} = \lambda_\theta = \lambda_{\hat{\mathbf{i}}} = \lambda_{\perp\hat{\omega}}$  by a new 1-vector  $\hat{\omega}$ ,  $|\hat{\omega}|=1$ , which is the normal vector to the rotation plane  $\hat{\omega} \perp \lambda_{\mathbf{U}}$ . This *pqg-1*-vector subject is perpendicular to the 'paper' plane in Figure 5.46 and Figure 5.47 and indicates the **direction**, that light has to face the observer from the rotation plane. – Imagine that you look perpendicular to the figure plane surface right on.<sup>260</sup>

5.4.5.5. The Duality of Direction

The light you as the reader of this book receive and see has two characteristics:

- You receive it as a particle expressed as a momentum *pqg-1*-vector  $\frac{\hbar}{c} \omega \hat{\omega}$ , or
- You receive it as a transversal plane wave expressed through a unitary development 1-rotor  $\mathbf{U}_\omega(t) = e^{\pm i\theta} = e^{\pm i\omega t} \sim \mathbf{b}\mathbf{e} = \mathbf{b} \cdot \mathbf{e} + \mathbf{b} \wedge \mathbf{e}$ . We presume orthogonality  $\mathbf{b} \cdot \mathbf{e} = 0$ . Then the field bivector  $\mathbf{F} = \mathbf{b} \wedge \mathbf{e} \sim \|\hat{\mathbf{i}}$  gives the plane *pqg-2 direction* of the wavefront.

These two are mutual dual following the complementarity principle between particle *pqg-1 direction* and the plane wavefront *pqg-2 direction*, where the plane rotor is a combination of a scalar *pqg-0* quality without **direction** and the bivector plane **direction** *pqg-2 quality* as an even geometric algebra.

The idea of a particle picture of a physical **entity** will be concerned with the odd part of Geometric Algebra. More of this below in the following chapters.

<sup>259</sup> We will below in section 5.7.5 interpret the development as a **direction** unit called  $\gamma_0$  as a 1-vector in Space-Time-Algebra.

<sup>260</sup> You cannot see  $\hat{\omega}$  in perspective, which causes the viewing angle loses its meaning. (An abstraction automatic don in your brain.) Later in Section 5.7.4, this will give analytical meaning by the null direction of a Lorentz rotation. The duality concept problem will be further analysed in chapter 6.

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